

University College London
Department of Physics and Astronomy
2B21 Mathematical Methods in Physics & Astronomy
Suggested Solutions for Problem Sheet M7 (2003–2004)

1. We are given that

$$\int_{-1}^{+1} P_m(x)P_n(x) dx = \delta_{nm} .$$

Starting with the cases where $n \neq m$, since $P_1(x)$ is odd and $P_0(x)$ and $P_2(x)$ are even, it follows immediately that

$$\int_{-1}^{+1} P_0(x)P_1(x) dx = \int_{-1}^{+1} P_2(x)P_1(x) dx = 0 . \quad [1]$$

Furthermore,

$$\int_{-1}^{+1} P_2(x)P_0(x) dx = \frac{1}{2} \int_{-1}^{+1} (3x^2 - 1) dx = \frac{1}{2} [x^3 - x]_{-1}^{+1} = 0 . \quad [2]$$

For $n = m = 0$,

$$\int_{-1}^{+1} [P_0(x)]^2 dx = \int_{-1}^{+1} dx = 2 . \quad [1]$$

For $n = m = 1$,

$$\int_{-1}^{+1} [P_1(x)]^2 dx = \int_{-1}^{+1} x^2 dx = \frac{2}{3} . \quad [1]$$

For $n = m = 2$,

$$\int_{-1}^{+1} [P_2(x)]^2 dx = \frac{1}{4} \int_{-1}^{+1} (9x^4 - 6x^2 + 1) dx = \frac{1}{4} \left[\frac{9}{5}x^5 - 2x^3 + x \right]_{-1}^{+1} = \frac{2}{5} . \quad [1]$$

If $P_3(x) = a[x^3 + bx^2 + cx + d]$, orthogonality with $P_0(x)$ requires that

$$\int_{-1}^{+1} [x^3 + bx^2 + cx + d] dx = \frac{2}{3}b + 2d = 0 . \quad [1]$$

Similarly, orthogonality with $P_1(x)$ necessitates

$$\int_{-1}^{+1} [x^4 + bx^3 + cx^2 + dx] dx = \frac{2}{5} + \frac{2}{3}c = 0 . \quad [1]$$

Finally, imposing orthogonality with respect to $P_2(x)$ means that

$$\frac{3}{2} \int_{-1}^{+1} [x^5 + bx^4 + cx^3 + dx^2] dx + \frac{1}{2} \int_{-1}^{+1} [x^3 + bx^2 + cx + d] dx = 0 .$$

Now the second integral vanishes because of the orthogonality to $P_0(x)$. Hence

$$\frac{3}{5}b + \frac{1}{2}d = 0 . \quad [1]$$

This result is incompatible with the previous relation obtained between b and d . Hence $b = d = 0$. Students might get this result by claiming that $P_3(x)$ is an odd function. Though true, this was not given in the question and would only receive a maximum credit of **three** marks.

Using the result for c , the Legendre polynomial reduces to

$$P_3(x) = a \left[x^3 - \frac{3}{5}x \right] . \quad [1]$$

The easiest way of determining the value of a is from the condition that the Legendre polynomials are normalised by $P_n(x = 1) = 1$. Therefore

$$1 = a \left[1 - \frac{3}{5} \right] = \frac{2}{5}a ,$$

and $a = \frac{5}{2}$. [3]

Alternatively, For $n = m = 2$,

$$\begin{aligned} \int_{-1}^{+1} [P_3(x)]^2 dx &= a^2 \int_{-1}^{+1} \left(x^6 - \frac{6}{5}x^4 + \frac{9}{25}x^2 \right) dx \\ &= a^2 \left[\frac{1}{7}x^7 - \frac{6}{25}x^5 + \frac{3}{25}x^3 \right]_{-1}^{+1} = \frac{8}{175}a^2 = \frac{2}{7} . \end{aligned}$$

Hence $a^2 = 25/4$ and we find once more that $a = \frac{5}{2}$. [3]

2. Expanding the left hand side of

$$g(x, t) = \frac{\exp(-xt/(1-t))}{1-t} = \sum_{n=0}^{\infty} L_n(x) t^n$$

in powers of t gives

$$\begin{aligned} g(x, t) &= (1+t+t^2) \exp(-xt(1+t)) + O(t^3) \approx (1+t+t^2)(1-xt-xt^2+\frac{1}{2}x^2t^2) \\ &= 1+t(1-x) + t^2(\frac{1}{2}x^2-2x+1) + O(t^3) . \end{aligned} \quad [1]$$

Hence $L_0(x) = 1$, $L_1(x) = 1-x$, [1]

and $L_2(x) = \frac{1}{2}(x^2-4x+2)$. [1]

Differentiating the generating function with respect to x gives

$$\frac{\partial g(x, t)}{\partial x} = - \left(\frac{t}{1-t} \right) g(x, t) = \sum_{n=0}^{\infty} L'_n(x) t^n . \quad [1]$$

Expanding $g(x, t)$ as a power series and multiplying through by $(1-t)$ results in

$$t \sum_{n=0}^{\infty} L_n(x) t^n = (1-t) \sum_{n=0}^{\infty} L'_n(x) t^n .$$

$$\sum_{n=0}^{\infty} L_n(x) t^{n+1} = \sum_{n=0}^{\infty} L'_n(x) t^n - \sum_{n=0}^{\infty} L'_n(x) t^{n+1} . \quad [1]$$

Now change the summation index $n \rightarrow n+1$ in the first term on the right hand side and compare the coefficients of t^{n+1} ;

$$L_n(x) = L'_{n+1}(x) - L'_n(x) . \quad [2]$$

Now

$$g(x, t) = \frac{\exp(-xt/(1-t))}{1-t} = \sum_{n=0}^{\infty} L_n(x) t^n$$

$$g(x, u) = \frac{\exp(-xu/(1-u))}{1-u} = \sum_{n=0}^{\infty} L_n(x) u^n \quad [1]$$

$$I = \int_0^{\infty} e^{-x} g(x, t) g(x, u) dx = \frac{1}{(1-t)(1-u)} \int_0^{\infty} \exp \left[-x \left(1 + \frac{t}{1-t} + \frac{u}{1-u} \right) \right] dx$$

$$= \frac{1}{(1-t)(1-u)} \int_0^{\infty} \exp \left[-x \left(\frac{1-ut}{(1-t)(1-u)} \right) \right] dx = \frac{1}{1-ut} = \sum_{n=0}^{\infty} u^n t^n . \quad [2]$$

Using the expansion formulae, the integral must also be given by

$$I = \sum_{n=0}^{\infty} \sum_{m=0}^{\infty} t^n t^m \int_0^{\infty} e^{-x} L_n(x) L_m(x) dx . \quad [1]$$

On the left hand side the power of u is always equal to the power of t . Hence only $n = m$ is non-vanishing on the RHS so that the required integral is proportional to δ_{nm} . [1]

Then, comparing powers of $u^n t^n$ on both sides, this shows that

$$\int_0^{\infty} e^{-x} L_n(x) L_n(x) dx = 1 . \quad [1]$$

This gives the normalisation integral for the Coulomb wave functions that is of use in the 2B22 course.