## 9 Maxwell's Equations in differential form

## Summary of Maxwell's equations 9.1

We have now seen all four of Maxwell's equations. They are

$$\begin{aligned}
\int_{S} \underline{E} \cdot d\underline{a} &= \frac{Q}{\epsilon_{0}} \quad \text{Gauss' law} \\
\int_{S} \underline{B} \cdot d\underline{a} &= 0 \quad \text{Gauss' law in magnetism} \\
\oint \underline{E} \cdot d\underline{s} &= -\frac{d\Phi_{B}}{dt} \quad \text{Faraday's law} \\
\begin{pmatrix} 9.2 \\ 9.3 \end{pmatrix} \\
\int_{S} \underline{B} \cdot d\underline{s} &= -\frac{d\Phi_{B}}{dt} \quad \text{Faraday's law} \\
\begin{pmatrix} 9.3 \\ 9.3 \end{pmatrix} \\
\begin{pmatrix} 0.4 \\ 0.4 \end{pmatrix} \\
\end{pmatrix}$$

$$\int_{S} \underline{B} \cdot d\underline{a} = 0 \quad \text{Gauss' law in magnetism}$$
(9.2)  
$$\oint \underline{E} \cdot d\underline{s} = -\frac{d\Phi_{B}}{dt} \quad \text{Faraday's law}$$
(9.3)

$$\oint \underline{B} \cdot d\underline{s} = \mu_0 I + \epsilon_0 \mu_0 \frac{d\Phi_E}{dt} \quad \text{Ampere - Maxwell law} \quad (9.4)$$

We have also seen their differential forms:

$$\left( \begin{array}{cccc} \underline{\nabla} \cdot \underline{E} &=& \frac{\rho}{\epsilon_0} & \text{Gauss' law} & (9.5) \\ \underline{\nabla} \cdot \underline{B} &=& 0 & \text{Gauss' law in magnetism} & (9.6) \\ \underline{\nabla} \times \underline{E} &=& -\frac{\partial \underline{B}}{\partial t} & \text{Faraday's law} & (9.7) \\ \underline{\nabla} \times \underline{B} &=& \mu_0 \underline{J} + \epsilon_0 \mu_0 \frac{\partial \underline{E}}{\partial t} & \text{Ampere - Maxwell law} & (9.8) \end{array} \right)$$

## Waves in free space 9.2

To summarise, Maxwell's equations in differential form in free space (no charges or currents present) are

$$\underline{\nabla} \cdot \underline{E} = 0 \tag{9.9}$$

$$\underline{\nabla} \cdot \underline{B} = 0 \tag{9.10}$$

$$\underline{\nabla} \times \underline{\underline{E}} = -\frac{\partial \underline{\underline{B}}}{\partial t} \tag{9.11}$$

$$\underline{\nabla} \times \underline{B} = \mu_0 \epsilon_0 \frac{\partial \underline{E}}{\partial t} \tag{9.12}$$

Now, taking the curl of 9.11 we get

$$\underline{\nabla} \times (\underline{\nabla} \times \underline{E}) = -\frac{\partial (\underline{\nabla} \times \underline{B})}{\partial t} = -\mu_0 \epsilon_0 \frac{\partial^2 \underline{E}}{\partial t^2}$$
(9.13)

where we have used equation 9.12.

We also have a general result for  $\underline{\nabla} \times (\underline{\nabla} \times \underline{B})$  which gives

$$\underline{\nabla} \times (\underline{\nabla} \times \underline{E}) = \underline{\nabla} (\underline{\nabla} \cdot \underline{E}) - (\underline{\nabla} \cdot \underline{\nabla}) \underline{E}$$
(9.14)

But  $\underline{\nabla} \cdot \underline{E} = 0$  in empty space, so

$$\underline{\nabla} \times (\underline{\nabla} \times \underline{E}) = -\nabla^2 \underline{E} \tag{9.15}$$

and hence

$$\left(\nabla^2 \underline{E} = \mu_0 \epsilon_0 \frac{\partial^2 \underline{E}}{\partial t^2} \qquad (9.16)\right)$$

 $\nabla^2$  is the Laplacian operator

$$rac{\partial^2}{\partial x^2}+rac{\partial^2}{\partial y^2}+rac{\partial^2}{\partial z^2}$$

We can, by the same approach, obtain a similar equation for  $\underline{B}$ 

$$\left(\nabla^2 \underline{B} = \mu_0 \epsilon_0 \frac{\partial^2 \underline{B}}{\partial t^2} \qquad (9.17)\right)$$

Equations 9.16 and 9.17 are three-dimensional wave equations. The general 3-D wave equation is

$$\nabla^2 \underline{A} = \frac{1}{v^2} \frac{\partial^2 \underline{A}}{\partial t^2}$$

where v is the speed of propagation of the wave. By comparing with 9.16 and 9.17 we see that the oscillations in  $\underline{E}$  and  $\underline{B}$  propagate at a speed

$$c = \frac{1}{\sqrt{\mu_0 \epsilon_0}} \tag{9.18}$$

which is the speed of light. Equations 9.16 and 9.17 describe waves comprising oscillations in electric and magnetic fields that propagate at the speed of light in a vacuum, in the absence of electric charges or currents.