

A G Polnarev. Mathematical aspects of cosmology (MAS347), 2008. IV. Relativistic Cosmological Models, Lecture 25.

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LECTURE 25

25.1. The three-dimensional space of constant curvature

According to the cosmological principle the Universe is the same everywhere, as a consequence, The three-dimensional space is curved in the same way everywhere, which means that at each moment of time the metric of the space is the same at all points. To obtain such a metric let us start from the following geometrical analogy. Let us consider the two-dimensional sphere in the flat three-dimensional space. In this case the element of length is

$$dl^2 = (dx^1)^2 + (dx^2)^2 + (dx^3)^2.$$

The equation of a sphere of radius a in the three-dimensional space has the form

$$(x^1)^2 + (x^2)^2 + (x^3)^2 = a^2.$$

The element of length on the two-dimensional sphere can be obtained if one expresses dx^3 in terms of dx^1 and dx^2 . From the equation for sphere we have

$$x^1 dx^1 + x^2 dx^2 + x^3 dx^3 = 0,$$

hence

$$dx^3 = -\frac{x^1 dx^1 + x^2 dx^2}{x^3} = -\frac{x^1 dx^1 + x^2 dx^2}{\sqrt{a^2 - (x^1)^2 - (x^2)^2}}.$$

Substituting dx^3 into dl we have

$$dl^2 = (dx^1)^2 + (dx^2)^2 + \frac{(x^1 dx^1 + x^2 dx^2)^2}{a^2 - (x^1)^2 - (x^2)^2}.$$

Introducing the "polar" coordinates instead of x^1 and x^2

$$x^1 = r \cos \phi,$$

$$x^2 = r \sin \phi,$$

we obtain

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$$\begin{aligned} dl^2 &= (dr \cos \phi - r \sin \phi d\phi)^2 + (dr \sin \phi + r \cos \phi d\phi)^2 + \\ &\frac{r \cos \phi (dr \cos \phi - r \sin \phi d\phi) + r \sin \phi (dr \sin \phi + r \cos \phi d\phi)}{a^2 - r^2} = \frac{a^2 dr^2}{a^2 - r^2} + r^2 d\phi^2 = \\ &= \frac{dr^2}{1 - \frac{r^2}{a^2}} + r^2 d\phi^2. \end{aligned}$$

Now we can repeat step by step the previous derivation, by considering the geometry of the three-dimensional space as the geometry on the three-dimensional hypersphere in some fictitious four-dimensional space (don't confuse with the physical four-dimensional space-time). In this case the element of length is

$$dl^2 = (dx^1)^2 + (dx^2)^2 + (dx^3)^2 + (dx^4)^2.$$

The equation of a sphere of radius a in the four-dimensional space has the form

$$(x^1)^2 + (x^2)^2 + (x^3)^2 + (x^4)^2 = a^2.$$

The element of length on the three-dimensional hypersphere, which represents the three-dimensional space of constant curvature, can be obtained, if one expresses dx^4 in terms of dx^1 , dx^2 and dx^3 . From the equation for hypersphere we have

$$x^1 dx^1 + x^2 dx^2 + x^3 dx^3 + x^4 dx^4 = 0,$$

hence

$$dx^4 = -\frac{x^1 dx^1 + x^2 dx^2 + x^3 dx^3}{x^4} = -\frac{x^1 dx^1 + x^2 dx^2 + x^3 dx^3}{\sqrt{a^2 - (x^1)^2 - (x^2)^2 - (x^3)^2}}.$$

Substituting dx^4 into dl we have

$$dl^2 = (dx^1)^2 + (dx^2)^2 + (dx^3)^2 + \frac{(x^1 dx^1 + x^2 dx^2 + x^3 dx^3)^2}{a^2 - (x^1)^2 - (x^2)^2 - (x^3)^2}.$$

Introducing the "spherical" coordinates instead of x^1 , x^2 and x^3

$$\begin{aligned} x^1 &= r \sin \theta \cos \phi, \\ x^2 &= r \sin \theta \sin \phi, \\ x^3 &= r \cos \theta, \end{aligned}$$

we obtain

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$$\begin{aligned}
 dl^2 &= (dr \sin \theta \cos \phi + r \cos \theta d\theta - r \sin \theta \sin \phi d\phi)^2 + (dr \sin \theta \sin \phi + r \cos \theta \cos \phi d\theta + r \sin \theta \sin \phi d\phi)^2 + \\
 &\quad + (dr \cos \theta - r \sin \theta d\theta)^2 + \\
 &\quad \frac{1}{a^2 - r^2} [r \sin \theta \cos \phi (dr \sin \theta \cos \phi + r \cos \theta d\theta - r \sin \theta \sin \phi d\phi) + \\
 &\quad + r \sin \theta \sin \phi (dr \sin \theta \sin \phi + r \cos \theta \cos \phi d\theta + r \sin \theta \sin \phi d\phi) + r \cos \theta (dr \cos \theta - r \sin \theta d\theta)] = \\
 &= \frac{a^2 dr^2}{a^2 - r^2} + r^2 (d\theta^2 + \sin^2 \theta d\phi^2) = \\
 &= \frac{dr^2}{1 - \frac{r^2}{a^2}} + r^2 (d\theta^2 + \sin^2 \theta d\phi^2).
 \end{aligned}$$

Taking into account that $r \leq a$ we can introduce instead of r the new lagrangian radial coordinate χ such that

$$r = a \sin \chi \quad \text{and} \quad dr = a \cos \chi d\chi,$$

as a result dl can be rewritten as

$$dl^2 = a^2 [d\chi^2 + \sin^2 \chi (d\theta^2 + \sin^2 \theta d\phi^2)].$$

Now we can write the metric interval for the four dimensional space time as

$$ds^2 = c^2 dt^2 - a^2 [d\chi^2 + \sin^2 \chi (d\theta^2 + \sin^2 \theta d\phi^2)].$$

We can repeat all these calculation for the three-dimensional space of the negative constant curvature for that one should replace the equation for the hypersphere by

$$(x^1)^2 + (x^2)^2 + (x^3)^2 - (x^4)^2 = a^2$$

[see cw5. Obviously, when $a \rightarrow \infty$ we obtain the case of the spatially flat space].

This method is called the method of embedding diagrams. The Geometry on different surfaces of constant curvature is shown in **Fig.12.2**. We will see later then in the relativistic cosmology the curvature of the three-dimensional space is related with the density parameter as it is shown on this figure.

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The fact that the Universe is expanding means that instead of constant a we should introduce a scale factor $R(t)$ and finally we obtain the famous the Friedmann-Lematre-Robertson-Walker metric for expanding Universe

25.2. Friedmann-Lematre-Robertson-Walker metric

$$ds^2 = c^2 dt^2 - R^2(t)[d\chi^2 + f^2(\chi)(d\theta^2 + \sin^2 \theta d\phi^2)],$$

where

$$f = \left\{ \begin{array}{ll} \sin \chi & \text{for constant positive curvatur} \\ \sinh \chi & \text{for constant negative curvatur} \\ \chi & \text{for zero curvatur} \end{array} \right\}.$$

It is convenient to introduce another time coordinate, η , defined by

$$cdt = ad\eta.$$

Then

$$ds^2 = R(\eta)[d\eta^2 - d\chi^2 - f^2(\chi)(d\theta^2 + \sin^2 \theta d\phi^2)].$$

Now using the EFEs we can obtain required equations for $R(t)$, in other words we can obtain relativistic cosmological models based on the EFEs.