accuracy of hydrostatic assumption **±**▲**444▶9**? **c**

In deriving the <u>equation of hydrostatic support</u>, it has been assumed that the gravity and pressure forces are balanced in a star. How valid an assumption is this?

Let us consider again the element of mass in <u>figure 6</u>. We have seen that the outward force acting on the element is given by

$P_{\rm r}\delta S$

and the inward force acting on the element is given by

$$P_{r+\delta r} \delta S + (GM_r / r^2) \rho_r \delta S \delta r$$

If the inward and outward forces are not equal, there will be a resultant force acting on the element which will give rise to an acceleration, *a*. This resultant force is given by $\rho_r \,\delta S \,\delta r \,a$, where we have simply multiplied the acceleration of the element by its mass. Hence we can write

$$P_{r+\delta r}\,\delta S + (GM_r\,/\,r^2)\,\rho_r\,\delta S\,\delta r - P_r\,\delta S = \rho_r\,\delta S\,\delta r\,a,$$

where we have assumed an inward-acting resultant force. If we are considering an infinitesimal element, we can write,

$$P_{r+\delta r} = P_r + (dP_r / dr) \delta r$$
 (in the limit $\delta r \rightarrow 0$, $\delta P_r / \delta r = dP_r / dr$, where $\delta P_r = P_{r+\delta r} - P_r$).

Combining these last two equations and rearranging, we obtain:

$$dP_r / dr + (GM_r / r^2) \rho_r = \rho_r a$$

Recalling that the acceleration due to gravity, g_r is given by GM_r / r^2 , and, for the sake of brevity, dropping the subscript r_r we may write:

 $dP/dr + \rho g = \rho a.$

This is a generalised form of the <u>equation of hydrostatic support</u>. We are now in a position to determine what happens if there is a resultant force acting on the element, i.e. if the sum of the two terms on the left-hand side of the above equation is not zero. Suppose that their sum is a small fraction β of the gravitational term, i.e.

$\beta \rho g = \rho a.$

This means that there will be an inward acceleration given by

 $a = \beta g.$

If the element starts from rest with this acceleration, its inward displacement *s* after a time *t* will be given by:

$$s = \frac{\gamma_2 a t^2}{2} = \frac{\gamma_2 \beta g t^2}{2}.$$

Rewriting the above equation in terms of *t*, we obtain:

$$t = (1 / \beta^{\gamma_2}) \times (2s / g)^{\gamma_2}.$$

If we allow the element to fall all the way to the centre of the star, we can replace *s* in the above equation by *r* and then subsitute $g = GM / r^2$, giving:

$$t = (1 / \beta^{\gamma_2}) \times (2r^3 / GM)^{\gamma_2}.$$

The last term is known as the <u>dynamical timescale</u>, t_d , i.e.

$$t_{\rm d} = (2r^3 / GM)^{\frac{1}{2}},$$

so we can write:

$$t = t_{\rm d} / \beta^{\gamma_2}.$$

This equation gives the time it would take a star to collapse if the forces are out of balance by a factor β .

Fossil and geological records indicate that the properties of the Sun have not changed significantly for at least 10^9 years (3 × 10^{16} s) and we know that the dynamical timescale for the Sun is approximately 2000 s (calculated by substituting values for the mass and radius of the Sun in the above expression for t_d). Hence, in the case of the Sun, we find that β can be no greater than 10^{-27} .

Not all stars are like the Sun, of course. During their lives, all stars undergo periods of radial expansion and/or contraction and at these times β will be much greater than this value. In such cases, the generalised form of the equation of hydrostatic support derived above must be used. Nevertheless, most stars *are* like the Sun (i.e. on the main sequence) and so we may conclude that:

the equation of hydrostatic support must be true to a very high degree of accuracy.

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