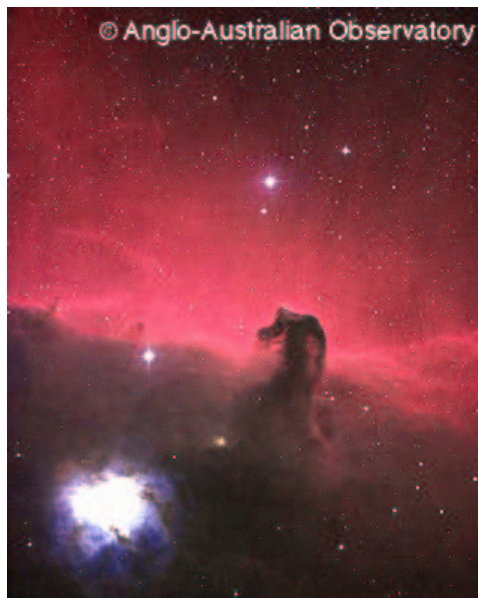


Star Formation (I)



Orion Nebula



STAR FORMATION (ZG: 15.3; CO: 12)

Star-Forming Regions

a) Massive stars

- born in **OB associations** in warm molecular clouds
- produce **brilliant HII regions**
- shape their environment
 - ▷ photoionization
 - ▷ stellar winds
 - ▷ supernovae
- induce further (low-mass) star formation?

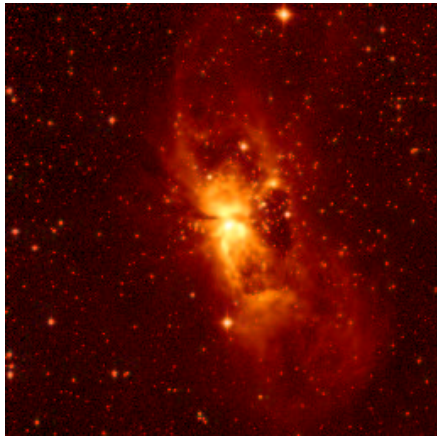
b) Low-mass stars

- born in **cold, dark molecular clouds** ($T \simeq 10$ K)
- Bok globules
- near massive stars?
- **recent:** most low-mass stars appear to be born in **cluster-like environments**
- **but:** most low-mass stars are not found in clusters → **embedded clusters do not survive**

Relationship between massive and low-mass star formation?

- ▷ massive stars trigger low-mass star formation?
- ▷ massive stars terminate low-mass star formation?

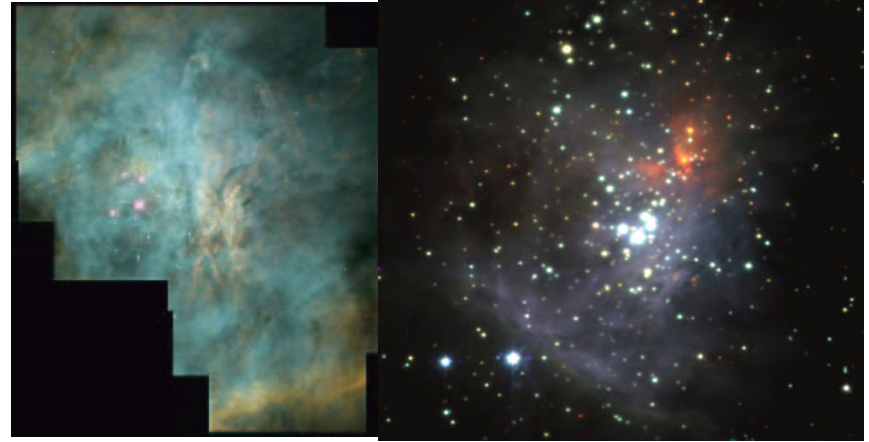
Star Formation (II)



S 106

massive star +
cluster of low-mas stars

Star Formation (III)

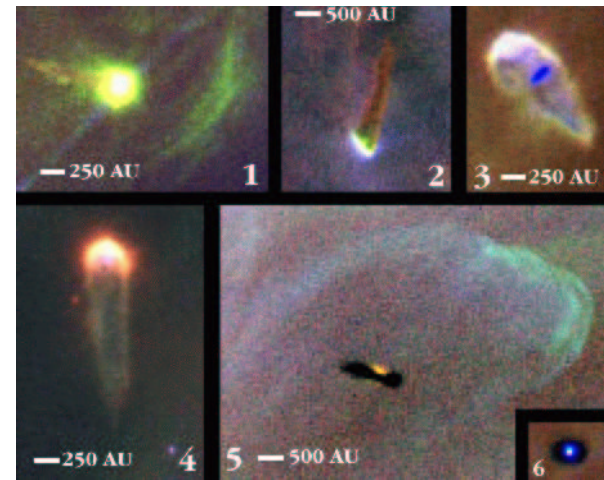


The Trapezium Cluster (IR)

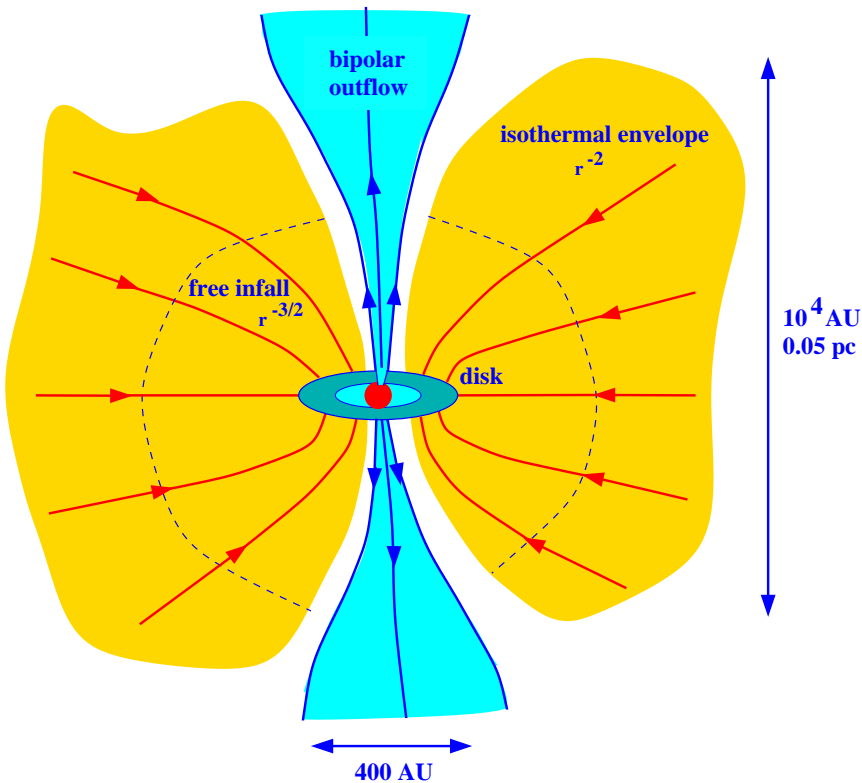


Bok globules

HST



Dusty Disks in Orion (seen as dark silhouettes)



Protostar Structure

The Jeans Mass

- cool, **molecular cores** (H_2) can collapse when their mass exceeds the **Jeans Mass**

▷ no thermal pressure support if $P_c = \rho / (\mu m_H) kT < GM^2 / (4\pi R^4)$

▷ or $M > M_J \simeq 6 M_\odot \left(\frac{T}{10 \text{ K}} \right)^{3/2} \left(\frac{n_{\text{H}_2}}{10^{10} \text{ m}^{-3}} \right)^{-1/2}$

What triggers star formation?

- **observed molecular clouds** often have masses \gg Jeans mass
- **but:** no evidence for large-scale collapse

→ support required

▷ cannot be thermal (**Jeans mass!** $v_{\text{th}} \ll v_{\text{virial}}$)

▷ **supersonic turbulence:** possible, but: rapid shock dissipation

▷ **magnetic fields:** requires $\rho v_{\text{virial}}^2 \sim B^2 / 2\mu_0 \rightarrow B \sim 1 - 10 \text{ nT}$ (o.k. consistent with observations)

- stars can form in regions that lose **magnetic support**
- **collisions** of cores (compression reduces Jeans mass)
- **compression** by nearby supernovae

Stellar Collapse

- **inside-out isothermal collapse** (i.e. efficient radiation of energy) from $\sim 10^6 R_\odot$ to $\sim 5 R_\odot$ (note this decreases the Jeans mass and possibly allows further fragmentation of the core)
- **timescale:** $t_{\text{dyn}} \sim 1/\sqrt{4G\rho} \sim 10^5 - 10^6 \text{ yr}$
- collapse **stops** when material becomes **optically thick** and can no longer remain isothermal (**protostar**)
- **central accretion rate:** \dot{M}

▷ hydrostatic equilibrium of an **isothermal sphere:**

$$c_s^2 = \frac{kT}{\mu m_H} = \frac{GM(r)}{r},$$

where c_s is the sound speed of the material, $M(r)$ the mass enclosed in radius r .

▷ $c_s = \text{constant}$ implies $M(r) \propto r$

→ for the density $\rho(r) = \frac{M_0}{4\pi r^2 R_0} = \frac{c_s^2}{4\pi r^2 G},$

where M_0 and R_0 are the total mass and total radius of the collapsing core.

▷ at radius r : mass-inflow rate \dot{M} is given by $\dot{M} = 4\pi r^2 \rho c_s$ (inflow velocity = sound speed)

▷ combining these equations, one obtains for the central accretion rate

$$\dot{M} = \frac{c_s^3}{G} = 2 \times 10^{-6} M_\odot \text{ yr}^{-1} \left(\frac{T}{10 \text{ K}} \right)^{3/2},$$

where $\mu = 2$ (molecular hydrogen) and

$$c_s = 0.2 \text{ km s}^{-1} \left(\frac{T}{10 \text{ K}} \right)^{1/2}.$$

▷ **note:** \dot{M} depends strongly on T , which in turn depends on the **cooling mechanisms** (CO molecules, dust, H_2 , etc.) and is dependent on the environment and metallicity.

- **the angular-momentum problem**

▷ each molecular core has a small amount of angular momentum (due to the velocity shear caused by the Galactic rotation)

▷ characteristic $\Delta v / \Delta R \sim 0.3 \text{ km/s/ly}$

→ characteristic, specific angular momentum

$$j \sim (\Delta v / \Delta R R_{\text{cloud}}) R_{\text{cloud}} \sim 3 \times 10^{16} \text{ m}^2 \text{ s}^{-1}$$

▷ cores cannot collapse directly

→ formation of an **accretion disk**

▷ characteristic disk size from angular-momentum conservation $j = r v_\perp = r v_{\text{Kepler}} = \sqrt{GM} r$

→ $r_{\text{min}} = j^2 / GM \sim 10^4 R_\odot \simeq 50 \text{ AU}$

- **Solution:** Formation of **binary systems and planetary systems** which store the angular momentum (**Jupiter:** 99 % of angular momentum in solar system)

→ **most stars should have planetary systems and/or stellar companions**

→ stars are initially **rotating rapidly** (spin-down for stars like the Sun by magnetic braking)

- **inflow/outflow:** $\sim 1/3$ of material accreted is ejected from the accreting protostar → bipolar jets

- **the magnetic field problem**

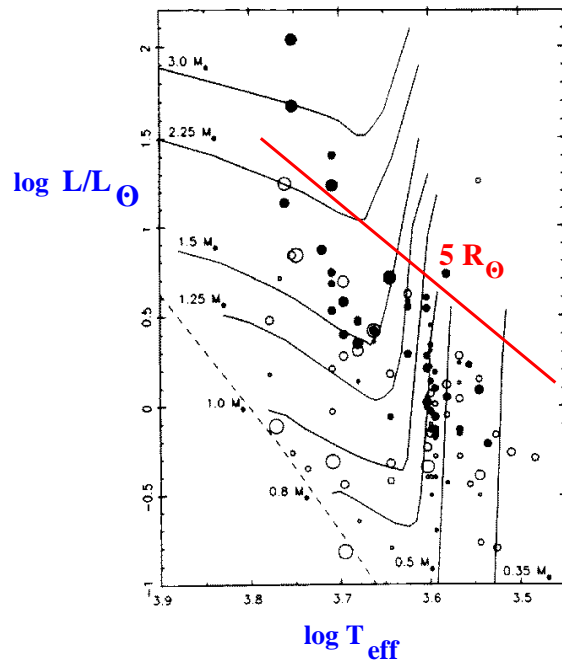
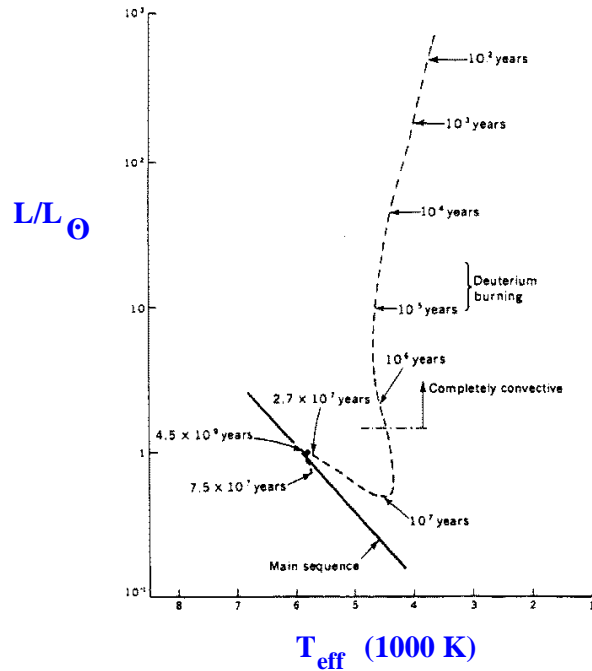
▷ using **magnetic flux conservation**

$$B(\text{star}) = B(\text{cloud}) (R_{\text{cloud}} / R_{\text{star}})^2 \sim 10^3 - 10^4 \text{ T (!)},$$

many order larger than observed

▷ efficient loss of magnetic field, perhaps related to bipolar jets

Pre-Main-Sequence Evolution

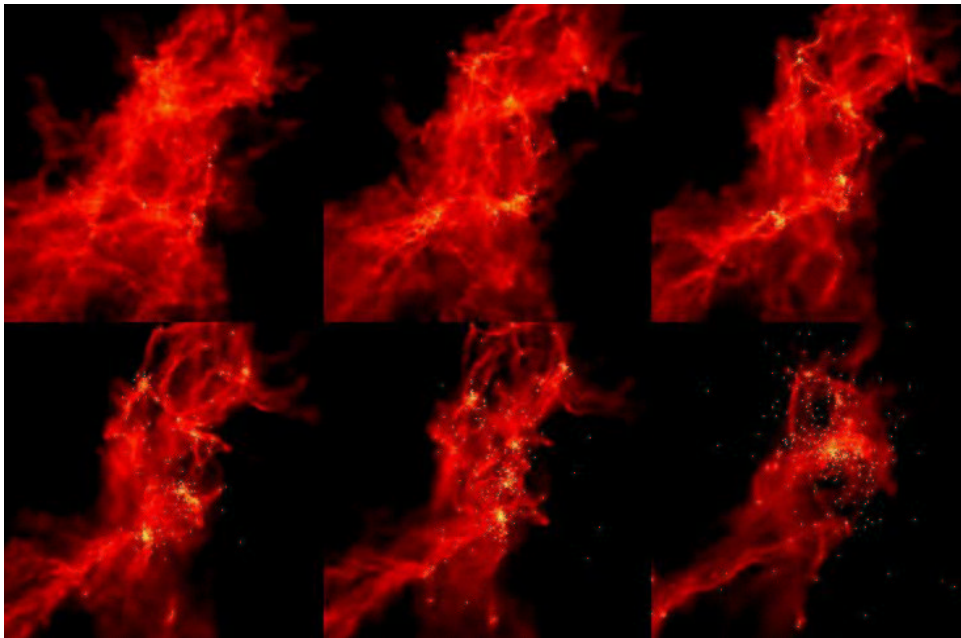


Pre-main-sequence evolution

- **Old picture:** stars are born with **large radii** ($\sim 100 R_{\odot}$) and slowly contract to the main sequence
 - ▷ energy source: **gravitational energy**
 - ▷ contraction stops when the central temperature reaches 10^7 K and H-burning starts (main sequence)
 - ▷ note: D already burns at $T_c \sim 10^6$ K \rightarrow temporarily halts contraction
- **Modern picture:** stars are born with **small radii** ($\sim 5 R_{\odot}$) and small masses
 - \rightarrow first appearance in the H-R diagram on the **stellar birthline** (where accretion timescale is comparable to Kelvin-Helmholtz timescale: $t_{\dot{M}} \equiv M/\dot{M} \sim t_{KH} = GM^2/(2RL)$)
 - ▷ continued accretion as **embedded protostars/T Tauri stars** until the mass is exhausted or accretion stops because of dynamical interactions with other cores/stars

Dynamical Star Formation

- stars generally do not seem to form in isolation, but in dense clusters
- simulation (Bonnell): $10^3 M_{\odot}$ cloud with radius 0.5 pc
 - collapse and fragmentation lead to the formation of ~ 400 stars in $\sim 0.5 \times 10^6$ yr with broad mass spectrum (but no magnetic fields considered in setting the initial conditions!)



- protostars form in collapsing cores ($R \sim 10^6 R_{\odot}$) and accrete from their cores at $\dot{M} \sim 2 \times 10^{-6} M_{\odot} \text{yr}^{-1}$ till the envelopes are disturbed by a collision with another core/star

▷ collision time: $t_{\text{coll}} \simeq 1/\sigma n v$

▷ where the collision cross section is given by the size of the core: $\sigma = \pi * (10^6 R_{\odot})^2$,

▷ the number density of colliding objects by $n \sim 10^3 / [(4\pi/3) \times (0.5 \text{ pc})^3]$ and

▷ the characteristic velocity by the dynamics of the cloud $v \sim \sqrt{GM/R} \simeq 3 \text{ km s}^{-1}$.

→ $t_{\text{coll}} \simeq 10^5 \text{ yr} \rightarrow M_{\text{star}} \sim \dot{M} \times t_{\text{coll}} \sim 10 M_{\odot}$

→ a collisional origin of the initial mass function?

The First Stars

- differences at zero metallicity:
 - ▷ no dust, no CO → higher T of star-forming cloud
 - larger Jeans mass → form very massive stars only?
- at $Z = 0$: very different stellar evolution (no CNO cycle) → different supernovae? Claim: pair-instability supernova: complete disruption of star in an energetic supernova (sometimes, also referred to as hypernova, not to be confused with GRB-related hypernova)
- but: observed nucleosynthesis from Pop III stars is not consistent with pair-instability supernovae
- formation of intermediate-mass black holes?
- Problem: it is not clear whether Pop III stars really should have existed as a significant population