

## WRG-11

## Energy in Wave Motion

Since we would like to consider waves in general, including harmonic waves such as  $y(x, t) = A \sin(kx - \omega t)$  which have an infinite extension in  $x$ , we shall consider energy per distance  $\lambda$ . That is we are considering the energy in one full oscillation.

For a section of a string between  $x$  and  $x + dx$ , its mass is approximately  $\mu dx$

$$\mu \equiv \frac{\text{mass}}{\text{unit length}}$$

Therefore the **kinetic energy** is  $dK = \frac{1}{2} \mu dx \left( \frac{\partial y}{\partial t} \right)^2$

and the kinetic energy per unit length is  $\frac{dK}{dx} = \frac{1}{2} \mu \left( \frac{\partial y}{\partial t} \right)^2$ . (1)

The **potential energy** is obtained by realising that the tension  $T$  acting on the string has extended it from  $dx$  to  $ds$ .

Hence the potential energy stored in segment  $dx$  is  $dU = T(ds - dx) = T dx \left( \frac{ds}{dx} - 1 \right)$ .

From the above triangle we have  $ds^2 = (dx)^2 + (dy)^2 = (dx)^2 \left( 1 + \left( \frac{\partial y}{\partial x} \right)^2 \right)$ .

Hence  $ds = dx \left( 1 + \left( \frac{\partial y}{\partial x} \right)^2 \right)^{\frac{1}{2}} \approx dx \left( 1 + \frac{1}{2} \left( \frac{\partial y}{\partial x} \right)^2 \right)$

and so  $\left( \frac{ds}{dx} - 1 \right) = \frac{1}{2} \left( \frac{\partial y}{\partial x} \right)^2$ .

Therefore  $\frac{dU}{dx} \approx \frac{1}{2} T \left( \frac{\partial y}{\partial x} \right)^2 = \frac{1}{2} \mu v^2 \left( \frac{\partial y}{\partial x} \right)^2$  as  $v^2 = \frac{T}{\mu}$ . (2)

As the wave equation must be obeyed, we can use the general solution  $y(x, t) = A \sin(kx - \omega t)$  in equations (1) and (2).

$$\text{So } \frac{\partial y}{\partial t} = -\omega A \cos(kx - \omega t) \quad \text{and so } \frac{dK}{dx} = \frac{1}{2} \mu \omega^2 A^2 \cos^2(kx - \omega t)$$

$$\frac{\partial y}{\partial x} = kA \cos(kx - \omega t) \quad \text{and so } \frac{dU}{dx} = \frac{1}{2} \mu k^2 v^2 A^2 \cos^2(kx - \omega t)$$

$$\text{But } kv = \omega, \quad \text{and so } \frac{dK}{dx} = \frac{dU}{dx}.$$

We see that the kinetic energy density and the potential energy density are equal.

$$\text{Now we wish to find the energy in one wavelength: } \int_0^\lambda \frac{dK}{dx} dx$$

$$K_\lambda = U_\lambda = \int_0^\lambda \frac{dK}{dx} dx = \frac{1}{2} \mu \omega^2 A^2 \int_0^\lambda \cos^2(kx - \omega t) dx \quad t \text{ is constant.}$$

To calculate the integral, change the variable to  $\theta = kx - \omega t$ .

Then  $d\theta = k dx$ .

$$\begin{array}{ll} \text{The limits are now} & x = 0, \quad \theta = -\omega t \\ & x = \lambda, \quad \theta = k\lambda - \omega t = 2\pi - \omega t \end{array}$$

So we have

$$\begin{aligned} \int_0^\lambda \cos^2(kx - \omega t) dx &= \frac{1}{k} \int_{-\omega t}^{2\pi - \omega t} \cos^2(\theta) d\theta = \frac{1}{2k} \int_{-\omega t}^{2\pi - \omega t} (1 + \cos 2\theta) d\theta \\ &= \frac{1}{2k} \left[ \theta + \frac{1}{2} \sin 2\theta \right]_{-\omega t}^{2\pi - \omega t} = \frac{1}{2k} \left[ 2\pi + \frac{1}{2} (\sin(4\pi - 2\omega t) - \sin(-2\omega t)) \right] \\ &= \frac{\pi}{k} \end{aligned}$$

$$\text{Therefore we have } K_\lambda = U_\lambda = \frac{\pi}{2} \frac{\mu \omega^2 A^2}{k} = \frac{1}{4} \mu \lambda \omega^2 A^2 = \pi^2 \mu \lambda f^2 A^2.$$

$$\text{The total energy is } E_\lambda = K_\lambda + U_\lambda = \frac{1}{2} \mu \lambda \omega^2 A^2.$$

Suppose we sit at a fixed point  $x$  and measure the total energy per second (power) passing that point, as the wave goes by. Since the wave is travelling at speed  $v$ , one wavelength  $\lambda$  takes a time  $\frac{\lambda}{v}$  sec to pass by.

Then power (joules per second) transported in the wave is

$$P = \frac{E_{\lambda}}{\left(\frac{\lambda}{v}\right)} = \frac{1}{2} \mu v \omega^2 A^2$$

Now, using  $v = \left(\frac{T}{\mu}\right)^{\frac{1}{2}}$  we define the quantity  $Z = \mu v = (T\mu)^{\frac{1}{2}}$  where  $Z$  is known as the **Impedance of the String**.

It is entirely a property of the string, while  $\omega^2 A^2$  refers to the nature of the propagating wave. The larger  $Z$ , the more power is needed to drive the wave of a given frequency and amplitude. Equivalently, for a given power input and frequency, a string of smaller impedance will have a correspondingly larger amplitude. The name comes from electricity where the impedance defines the *resistance* of a circuit for ac voltages.