## MATH 3305 General Relativity Problem sheet 9

Please hand in your solutions Friday, 18th December 2009
This is the last problem sheet and all problems can be done using the theory presented in class (as of December 11th).
Problem 1 (40 points) The Schwarzschild metric on $\mathbb{R}^{4}$ with coordinates $(t, r, \theta, \phi)$ is the metric tensor $g_{i j}$ with line element

$$
\begin{equation*}
d s^{2}=\left(1-\frac{r_{s}}{r}\right) d t^{2}-\frac{1}{c^{2}}\left(\left(1-\frac{r_{s}}{r}\right)^{-1} d r^{2}+r^{2} d \theta^{2}+r^{2} \sin ^{2} \theta d \phi^{2}\right) \tag{1}
\end{equation*}
$$

Here $r_{s} \geq 0$ is a constant.
(a) Write down the Lagrangian function for $d s^{2}$.
(b) Using the Lagrangian method, write down the geodesic equation for a curve

$$
(t(\lambda), r(\lambda), \theta(\lambda), \phi(\lambda)) .
$$

Problem 2 (20 points) Print out the Mathematica notebook SC.pdf on the course homepage. For each command $\operatorname{In}[1], \ldots, \operatorname{In}[14]$ give a short description of what it does, and show that the Einstein tensor of the metric given by equation (1) is zero.

Problem 3 ( 40 points) (Constant in Einstein's field equations) Suppose $g_{i j}$ is a Lorentz metric on a 4-manifold.
(a) Show that the Einstein field equations

$$
\begin{equation*}
R_{i j}-\frac{1}{2} R g_{i j}=\kappa T_{i j} \tag{2}
\end{equation*}
$$

can be rewritten as

$$
\begin{equation*}
R_{i j}=\kappa\left(T_{i j}-\frac{1}{2} g_{i j} T\right) \tag{3}
\end{equation*}
$$

where $T=T_{i j} g^{i j}$.
(b) Let us now assume that $\left(X^{0}=t, X^{1}=x, X^{2}=y, X^{3}=z\right)$ is an inertial frame for $M=\mathbb{R}^{4}$ and $\Phi(x, y, z)$ is a gravitational potential determined by a mass density $\varrho(x, y, z)$ so that

$$
\begin{equation*}
\nabla^{2} \Phi=4 \pi G \varrho \tag{4}
\end{equation*}
$$

Let $g_{i j}$ be the metric

$$
g_{i j}=\eta_{i j}-\left(\begin{array}{cccc}
-\frac{2 \Phi}{c^{2}} & & & \\
& -\frac{2 \Phi}{c^{4}} & & \\
& & -\frac{2 \Phi}{c^{4}} & \\
& & & -\frac{2 \Phi}{c^{4}}
\end{array}\right)
$$

We assume that products $\Phi^{2}, \Phi \frac{\partial \Phi}{\partial X^{r}}$ and $\frac{\partial \Phi}{\partial X^{r}} \frac{\partial \Phi}{\partial X^{s}}$ are so small that they can be neglected whenever they appear. In the lectures we then proved that geodesics of $g_{i j}$ coincide with the Newton equation of motion. Moreover, from the lectures

$$
\begin{aligned}
\Gamma_{00}^{i} & =\frac{\partial \Phi}{\partial X^{i}}, \quad \Gamma_{0 i}^{0}=\frac{1}{c^{2}} \frac{\partial \Phi}{\partial X^{i}}, \quad \Gamma_{j 0}^{i}=\Gamma_{i j}^{0}=\Gamma_{00}^{0}=0 \\
\Gamma_{j k}^{i} & =\frac{1}{c^{2}} \times \text { linear combination of }\left\{\frac{\partial \Phi}{\partial x}, \frac{\partial \Phi}{\partial y}, \frac{\partial \Phi}{\partial z}\right\}
\end{aligned}
$$

for all $i, j, k=1,2,3$. Using this show that

$$
R_{00}=\sum_{s=1}^{3} \frac{\partial \Gamma_{00}^{s}}{\partial X^{s}}=\nabla^{2} \Phi
$$

(c) Let $T_{i j}$ be the energy tensor for $\varrho$ given by

$$
T_{i j}=\left(\begin{array}{cccc}
c^{2} \varrho & & & \\
& 0 & & \\
& & 0 & \\
& & & 0
\end{array}\right)_{i j}
$$

Show that $T \approx c^{2} \varrho$.
(d) How should the constant $\kappa$ in the Einstein field equations be chosen such that equation (3) holds for $i=j=0$.

